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### EQUIVALENCE OF STAR-PRODUCTS ON SYMPLECTIC MANIFOLDS

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ABSTRACT. Besides properties related to the Hochschild cohomology of a symplectic manifold model in analytical dynamics with applications in quantum theory, this paper also shows the equivalence of two differential star-products, more specifically, that every differential star-product of two functions u and v on a symplectic manifold is equivalent to one whose linear term is half of the Poisson bracket of these functions, i.e.,  $\frac{1}{2}\{u, v\}$ .

#### 1. INTRODUCTION

Star-products were introduced in [2] to consider a particular deformation of the space  $C^{\infty}(M)$  of smooth functions on a symplectic manifold equipped with its double structure of associative algebra according to the usual pointwise multiplication of functions and the Poisson Lie algebra structure for a new approach of quantum mechanics. That is, a star-product is a formal deformation of these two algebraic structures. J. Vey [7] proved the existence of such deformations assuming that the third De Rham cohomology group of the manifold vanishes. Then, in 1983, M. De Wilde and P. B. A.

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Lecomte [4] proved the existence of a star-product on a symplectic manifold. In this paper, we will only be concerned with differential star-products, that is, star-products defined by a series of bidifferential operators on a symplectic manifold. A particular type of such star-products is defined.

#### 2. Preliminaries

**Definition 2.1.** A symplectic structure on a manifold M is a closed and nondegenerate 2-form  $\omega$  on M. The pair  $(M, \omega)$  is called symplectic manifold.

**Definition 2.2.** Let  $(M, \omega)$  be a symplectic manifold. A vector field X is said to be *locally Hamiltonian* if  $i_X \omega$  is an exact 1-form on M, where i is the interior product on M.

**Notation**. We shall denote by  $X_u$  the unique Hamiltonian vector field such that  $i_{X_u}\omega = du$  for  $u \in C^{\infty}(M)$ .

#### **Remarks:**

- (1) The space of symplectic vector fields modulo the space of Hamiltonian vector fields is isomorphic to the space of closed 1-forms modulo the space of exact 1-forms. That is, it is isomorphic to the first group of De Rham cohomology of M denoted by  $H^1(M, \mathbb{R})$ .
- (2) It is a consequence of the Poincare Lemma that every symplectic vector field is locally Hamiltonian.

**Definition 2.3.** Let  $(M, \omega)$  be a symplectic manifold, and let u and v be two smooth functions on M. The *Poisson bracket* of u and v, denoted by  $\{u, v\}$ , is defined by

$$\{u, v\} = X_u(v) = \omega(X_v, X_u).$$

Note that the Poisson bracket of u and v can also be defined in terms of the Lie derivative as follows:

$$\{u, v\} = -L_{X_u}(v) = L_{X_v}(u) = -i_{X_u} \circ i_{X_v}(\omega).$$

Defined in this way, the Poisson bracket can be viewed as a derivation on the space  $C^{\infty}(M)$  of smooth functions on the manifold M. The vector space  $C^{\infty}(M)$  equipped with the Poisson bracket is a

Lie algebra. The Poisson tensor  $\bigwedge$  is a 2-alternative vector field such that:

$$\{u, v\} = i \bigwedge (\mathrm{du} \wedge \mathrm{dv})$$

whose local co-ordinates  $\wedge^{ij}$  define a square matrix  $(\wedge^{ij})$  so that the inverse of this matrix is  $(\omega_{ij})$ , where  $\omega_{ij}$  are the components of the symplectic 2-form  $\omega$ . In what follows, we consider formal deformations of the associative structure of the algebra  $N = (C^{\infty}(M), \{,\})$ , that is, the deformations defined on the space  $N[\nu]$  of formal series in the formal parameter  $\nu$  with coefficients in N. The linear (or multilinear) map T on  $N[\nu]$  is formal if  $N[\nu] = \{f_{\nu}/f_{\nu} = \sum_{k=0}^{\infty} \nu^k f_k;$  $f_k \in N\}$ , and the linear (or multilinear) map T on  $N[\nu]$  is formal if  $T(\nu f_{\nu}) = \nu T(f_{\nu})$  (satisfied by each argument in the multilinear case).

#### 3. Star-product on a symplectic manifold $(M, \omega)$

Unless otherwise indicated, the manifold that we consider will be assumed paracompact. Furthermore, to set the following definition, we consider on  $C^{\infty}(M)$  the functions  $C_r$  (r a natural number) defined by  $C_r : C^{\infty}(M) \times C^{\infty}(M) \longrightarrow C^{\infty}(M)$  such that

- (1)  $C_r$  is a differential operator in both variables and annihilates constants;
- (2)  $C_0(u,v) = u \cdot v$  where  $\cdot$  is the pointwise multiplication of functions;
- (3)  $C_1(u,v) C_1(v,u) \equiv \{u,v\}$  is the Poisson bracket on  $(M,\omega)$ ;
- (4)  $C_k(u,v) = (-1)^k C_k(v,u)$  for  $k \ge 2$ ;
- (5)  $C_k(1, u) = C_k(u, 1) = 0.$

**Definition 3.1.** A star product (\*-product) on a symplectic manifold  $(M, \omega)$  is a formal deformation  $M_{\nu}$  defined on  $N[\nu]$  such that for all u and v in N,

$$M_{\nu}(u,v) = u \star v = u \star_{\nu} v = \sum_{r \ge 0} \nu^{r} C_{r}(u,v)$$

where  $C_r$  is defined as above and the constant function 1 is the unity. A  $\star$ -product is said to be of order k if  $u \star_{\nu} v = \sum_{r=0}^{k} C_r(u, v)$ .

The star-commutator on a symplectic manifold  $(M, \omega)$  is defined in terms of a star-product as follows

$$[u,v]_{\star} = u \star v - v \star u.$$

This is closely related to the commutator of observables in quantum theory. The symplectic manifold  $(M, \omega)$  is said to admit a deformation quantization provided a star-product always exists on M.

Note that the set  $N[\nu]$  is a Lie algebra whose adjoint representation is given by

$$ad_{\star}(u)(v) = [u, v]_{\star}$$
.

From the definition of the star-product, we have

$$[u, v]_{\star} = \nu\{u, v\} - 2\sum_{r=1}^{\infty} C_{2r+1}(u, v)\nu^{2r+1} + \dots$$

**Definition 3.2.** Two star-products,  $\star$  and  $\star'$ , on the symplectic manifold  $(M, \omega)$  are equivalent if there exists a series

$$T = \sum_{r=0}^{\infty} \nu^r T_r$$

where for each natural number r,  $T_r$  is a linear operator on N,  $T_0 = id_N$  is the identity function on N, and T is a linear bijection on  $N[\nu]$  which satisfies

$$T(u \star v) = T(u) \star' T(v).$$

The study of a star-product on its star-product algebra  $N[\nu] = C^{\infty}(M)[\nu]$  modulo this equivalence relation requires M. Gerstenhaber's deformation theory [5] which is based on the Hochschild cohomology of the algebra N.

**Definition 3.3.** A Hochschild p-*cochain* on the commutative algebra N is the p-linear map from  $N \times N \times ... \times N$  (p copies of N) into N, and a Hochschild coboundary operator on the algebra N is a map  $\partial$  satisfying the following property: for all p-cochains C on N,

$$(\partial C)(u_0, ..., u_p) = u_0 \star_{\nu} C(u_1, ..., u_p) + \sum_{r=1}^{p} (-1)^r C(u_0, ..., u_{r-1} \star_{\nu} u_r, ..., u_p) + (-1)^{p+1} C(u_0, ..., u_{p-1}) \star_{\nu} u_p .$$

**Remark**: In this paper, *p*-cochains, *p*-coboundaries, and the coboundary operator are related to Hochschild cohomology.

**Example.** A 1 or 2-coboundary  $\partial$  is given by

$$\begin{split} (\partial C_1)(u,v) &= uC_1(v) - C_1(u\cdot v) + C_1(u)\cdot v\\ (\partial C_2)(u,v,w) &= \\ & uC_2(v,w) - C_2(u\cdot v,w) + C_2(u,v\cdot w) - C_2(u,v)\cdot w\\ \text{where } C_1 \text{ and } C_2 \text{ are cochains.} \end{split}$$

**Definition 3.4.** A p-cochain C is said to be a *p*-cocycle if  $\partial C = 0$ , and it is said to be a coboundary if  $C = \partial B$  for some (p-1)-cochain B.

**Definition 3.5.** A p-cochain C is said to be

(1) differential if its variables are all differential operators;

(2) k-differential if each variable is a k-differential operator.

A p-cochain is said to annihilate constants if it vanishes for any constant among its variables.

It can be seen immediately that annihilating differential 1-cochains are cocycles, and that 1-cocycles are derivations on  $N = C^{\infty}(M)$ . In fact, from

$$0 = (\partial C)(u, v) = uC(v) - C(u \cdot v) + C(u) \cdot v,$$

one deduces that

$$C(u \cdot v) = uC(v) + C(u) \cdot v,$$

so that if the 1-cocycles define vector fields on M, then they define the 1-differential annihilating cochains.

**Definition 3.6.** The differential p-group of Hochschild cohomology of N is the quotient of the space of differential p-cocycles by the space of differential p-coboundaries.

From a p-cochain C and a q-cochain D, a (p + q)-cochain is defined by

 $(C \otimes D)(u_1, \dots, u_p, u_{p+1}, \dots, u_{p+q}) = C(u_1, \dots, u_p) + D(u_{p+1}, \dots, u_{p+q})$ 

and the coboundary operator is a graded derivation, that is,

$$\partial(C\otimes D) = (\partial C)\otimes D + (-1)^p C\otimes (\partial D).$$

When the cochain C is expressed in terms of local co-ordinates, its support supp C is the union of supports of coefficients.

**Proposition 3.1.** If C is a differential p-cocycle on  $C^{\infty}(\mathbb{R}^n)$ , then there exists a (p-1)-cochain B and a differential p-cocycle A such that  $C = \partial B + A$ . If C annihilates constants, then A and B can be chosen such that they also annihilate constants and their supports are contained in the support of C.

*Proof:* Let p = 1. Then every 1-cocycle is a vector field and so, the proposition holds trivially.

Now let us suppose that this result holds for a differential rcocycle  $C(u_1, ..., u_r)$  with  $1 \ge r \le (p-1)$  and  $u_1$  an operator of
order  $k \ (k > 1)$ . There exists a coboundary of first order as we can
let

$$C(u_1, ..., u_p) = \sum_{i_1, ..., i_k} \frac{\partial^k u_1}{\partial x_{i_1} ... \partial x_{i_k}} D_{i_1 ... i_k}(u_2, ..., u_p) + ...$$

where  $D_{i_1...i_k}$  are the symmetric cochains with respect to  $i_1, ..., i_k$ ; or, using the multi-index notation, it follows that

$$C = \sum_{|i|=k} \partial_i \otimes D_i + \dots$$

Thus, we can use the identities (1) and (2) above and the fact that  $\partial^2 = 0$  to get

$$\partial C = -\sum_{|i|=k} \partial_i \otimes \partial D_i + \dots$$

That is, if C is a *p*-cocycle, the coefficients of higher degree derivatives with respect to  $u_s$  (s = 1, ..., k) are (p - 1)-cocycles. By induction hypothesis,  $D_i = \partial E_i + F_i$ , where  $F_i$  are 1-differential and  $suppF_i \subset suppC$  for  $1 \leq i \leq p - 1$ . If

$$G = \sum_{|i|=k} \partial_i \otimes E_i + F_i \circ (\partial_i \otimes Id_{p-2}),$$

then a short calculation leads to

$$\partial G = -\sum_{|i|=k} \partial_i \otimes D_i + \dots$$

where ... stands for the terms in which the derivative of the first variable is of order less than k. Hence, the order of  $C + \partial G$  is at most k - 1 with respect to the first argument. By iteration, this order can be reduced to 1 with respect to the first argument. Suppose that

$$C = \sum_{i=1}^{n} \frac{\partial}{\partial x_i} \otimes D_i \; .$$

Then,

$$\partial C = -\sum_{i=1}^{n} \frac{\partial}{\partial x_i} \otimes \partial D_i,$$

and C is a cocycle if and only if  $D_i$  are cocycles, i.e.,

$$D_i = \partial E_i + F_i$$

where  $F_i$  are 1-differential and

$$C + \partial \sum_{i=1}^{n} \frac{\partial}{\partial x_i} \otimes E_i = \sum_{i=1}^{n} \frac{\partial}{\partial x_i} \otimes F_i.$$

**Proposition 3.2.** If C is a 1-differential p-cochain on  $\mathbb{R}^n$  whose alternating part is A, then

$$C = \partial B + A$$

where B is 2-differential and is determined by C, and  $suppB \subset suppC$ .

*Proof:* Let C be a 1-differential p-cochain annihilating constants. Then C has the following form:

$$C(u_1,...,u_p) = \sum_{i_1,i_2,...,i_p} C_{i_1i_2...i_p} \frac{\partial u_1}{\partial x_{i_1}} \cdots \frac{\partial u_p}{\partial x_{i_p}},$$

where the coefficients  $C_{i_1i_2...i_p}$  are given by

$$C_{i_1...i_p} = C(x_{i_1}, ..., x_{i_p}).$$

For a permutation  $\sigma$  on p elements  $\{1, ..., p\}$ , let

$$(\sigma C)(u_1, ..., u_p) = C(u_{\sigma^{-1}(1)}, ..., u_{\sigma^{-1}(p)})$$

define an action of the group  $\sigma \in \mathfrak{S}_p$  of all permutations of p elements on the p-cochains. For each permutation  $\sigma \in \mathfrak{S}_p$ , we define a 2-differential (p-1)-cochain  $\phi_{\sigma}(C)$  by

$$\partial \phi_{\sigma}(C) := C - \epsilon(\sigma)\sigma \cdot C$$

where  $\epsilon(\sigma)$  is the signature for  $\sigma \in \mathfrak{S}_p$ . In particular, if  $\tau$  is a transposition of consecutive integers, and if we consider a fixed *i* such that  $i \leq p-1$ , the (p-1)-cochain  $\phi_{\tau}(C)$ 

$$\phi_{\tau}(C)(u_1, ..., u_{p-1}) = (-1)^i \sum_{r,s} C(u_1, ..., u_{i-1}, x_r, x_s, u_{i+1}, ..., u_{p-1}) \frac{\partial^2 u_i}{\partial x_r \partial x_s}$$

satisfies

$$\partial \phi_{\tau}(C) = C + \tau \cdot C.$$

Then, we make two transpositions,  $\tau_1$  and  $\tau_2$ , of two consecutive integers and we have

$$\partial [\phi_{\tau_1}(\tau_2 C) - \phi_{\tau_2}(C)] = C - \tau_1 \tau_2 C.$$

For each  $\sigma \in \mathfrak{S}_p$ , we define a 2-differential (p-1)-cochain  $\phi_{\sigma}(C)$  by letting

$$\partial \phi_{\sigma}(C) := C - \epsilon(\sigma)\sigma \cdot C$$

where  $\epsilon(\sigma)$  is the signature of  $\sigma$ . Now, let

$$\phi(C) = \frac{1}{p!} \sum_{\sigma \in \mathfrak{S}_p} \phi_{\sigma}(C).$$

Then,

$$C = \partial \phi(C) + \frac{1}{p!} \sum_{\sigma \in \mathfrak{S}_p} \epsilon(\sigma) \sigma \cdot C$$

up to p!. Clearly, C is cohomologous to its antisymmetric part and that  $supp\phi(C) \subset suppC$ .

**Theorem 3.1.** A 1-differential p-cocycle C on a manifold M is the sum of a coboundary of a differential (p-1)-cochain B and a 1-differential skew-symmetric p-cocycle, A, i.e.,

$$C = \partial B + A.$$

If C annihilates constants, then B can also be chosen to annihilate constants.

*Proof:* Let  $\{U_{\lambda}\}_{\lambda \in \bigwedge}$  be a locally finite covering subordinated to a partition of unity  $\rho_{\lambda}$  of the manifold M. Then, a *p*-cocycle C can be written as a locally finite sum of *p*-cocycles

$$C = \sum_{\lambda \in \bigwedge} \rho_{\lambda} C_{\lambda}.$$

From propositions 3.1 and 3.2 above, we can write

$$\rho_{\lambda}C_{\lambda} = \partial B_{\lambda} + A_{\lambda}$$

where  $supp B_{\lambda} \subset U_{\lambda}$ . Let  $B = \sum_{\lambda \in \bigwedge} B_{\lambda}$  and  $A = \sum_{\lambda \in \bigwedge} A_{\lambda}$ , then  $C = \partial B + A$ , where A and B are locally finite, and are globally defined.  $\Box$ 

**Remark:** Note that in the symplectic case, the skew-symmetric 1-differential *p*-cocycle *A* can be written in terms of Hamiltonian vector fields as  $A(u_1, ..., u_p) = \alpha(X_{u_1}, ..., X_{u_p})$  where  $\alpha$  is therefore a one *p*-form and

$$C(u_1, ..., u_p) = (\partial B)(u_1, ..., u_p) + \alpha(X_{u_1}, ..., X_{u_p}).$$

**Definition 3.7.** A star-product,  $\star$ , on a symplectic manifold  $(M, \omega)$  is said to be differential if the 2-cochains  $C_{\tau}(u, v)$  defining it are bidifferential operators.

**Definition 3.8.** Two star-products,  $\star$  and  $\star'$ , on a symplectic manifold  $(M, \omega)$  are said to be differentially equivalent if there exists a series

$$T = \sum_{r=0}^{\infty} \nu^r T_r$$

where  $T_r$  are differential operators on  $N = C^{\infty}(M)$  and such that

$$T(u \star v) = T(u) \star' T(v).$$

Note that star-products that are differentially equivalent are equivalent. The following theorem asserts that when star-products are differential, they are differentially equivalent if and only if they are equivalent.

**Theorem 3.2.** Let  $\star$  and  $\star'$  be two differential star-products and let

$$T = \sum_{r \ge 0} \nu^r T_r,$$

where  $T_0 = id$  be an equivalence with  $T(u \star v) = T(u) \star' T(v)$ , then  $T_r$  are differential operators on N.

*Proof:* Let us suppose that the first k operators  $T_1, ..., T_k$  in T are differential and define  $T' = \sum_{r=0}^k \nu^r T_r$  and  $T'' = T'^{-1} \circ T$ . It is obvious that the form of T'' is

$$T''(u) = u + \nu^{k+1} T''_{k+1}(u) + \dots$$

Let us define a star-product  $\star''$  such that

$$u \star'' v = T^{-1}(T'(u) \star' T'(v))$$

and  $\star$  satisfies the following relation

$$u \star v = T'^{-1}(T''(u) \star'' T''(v)).$$

Then,  $T'' = T'^{-1} \circ T$  is an equivalence between two differential star-products,  $\star$  and  $\star''$ . If we consider only the (k+1)-th degree's terms in  $u \star v = T'^{-1}(T''(u) \star'' T''(v))$ , it follows that

$$(\partial T_{k+1}'')(u,v) = T_{k+1}''(u)v + uT_{k+1}''(v) - T_{k+1}''(u \cdot v)$$

is a 2-cocycle bidifferential and symmetric. Thus, from Theorem 3.1 above,  $\partial T_{k+1}''$  is the coboundary of a differential skew-symmetric 1-cochain plus a skew-symmetric differential 1-cocycle. The exact terms are symmetric so that the skew-symmetric part is equal to zero. Thus, there exists a differential 1-cochain *B* such that

$$\partial (T_{k+1}'' - B) = 0.$$

Hence,  $X = T''_{k+1} - B$  is a vector field on M, that is a derivation on  $N = C^{\infty}(M)$ . Thus,  $T''_{k+1} = B + X$  is differential. Now  $T_{k+1}$  is a linear combination of  $T_1, \ldots, T_k$  and  $T''_{k+1}$ , which are differential, and  $T_{k+1}$  is also. By induction, one concludes that T is differential.  $\Box$ 

The next proposition is an immediate application of Theorem 3.1 above.

**Proposition 3.3.** Every star-product is equivalent to a star-product such that its linear term in  $\nu$  is given by  $\frac{1}{2}\{u, v\}$ .

*Proof:* Let

$$u \star v = uv + \nu C_1(u, v) + \dots$$

be a star-product on a symplectic manifold  $(M, \omega)$ ; then, we need to show that  $C_1(u, v)$  is a Hochschild cocycle such that its skewsymmetric part is  $\frac{1}{2}\{u, v\}$ . That is, from Theorem 3.1, we should have

$$C_1(u,v) = uB(v) - B(u \cdot v) + B(u) \cdot v + \frac{1}{2} \{u,v\}$$

where B is a differential 1-cochain. Let  $T(u) = u + \nu B(u)$  and  $u \star' v = T(T^{-1}(u) \star T^{-1}(v))$ . Then

$$u \star' v = uv + \frac{1}{2}\nu\{u, v\} + \dots$$

Thus, T is a differential equivalence and so,  $\star'$  is a differential starproduct.

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